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Title: Combined effects of internal heat generation and buoyancy force on boundary layer over a vertical plate with a convective surface boundary condition

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Abstract: This paper considers the effect of buoyancy force and internal heat generation on laminar thermal boundary layer over a vertical plate with a convective surface boundary condition. We assumed that left surface of the plate is in contact with a hot fluid while a stream of cold fluid flows steadily over the right surface with a heat source that decays exponentially. Using a similarity variable, the steady state governing non-linear partial differential equations have been transformed into a set of coupled non-linear ordinary differential equations, which are solved numerically by applying shooting iteration technique together with fourth order Runge-Kutta integration scheme. The effects of Prandtl number, local Biot number, the internal heat generation parameter and the local Grashof number on the velocity and temperature profiles are illustrated and interpreted in physical terms. A comparison with previously published results on special case of the problem shows excellent agreement

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4 **Combined effects of internal heat generation and buoyancy force on boundary layer over a**
5 **vertical plate with a convective surface boundary condition**
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19 **Abstract**
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22 thermal boundary layer over a vertical plate with a convective surface boundary condition. We
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42 **Keywords:** Vertical plate; Convective boundary condition; Internal heat generation; Buoyancy
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49 **1. Introduction**
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51 Investigations of boundary layer flow and heat transfer of viscous fluids over a flat sheet are
52 important in many manufacturing processes, such as polymer extrusion, drawing of copper
53 wires, continuous stretching of plastic films and artificial fibers, hot rolling, wire drawing, glass-
54 fiber, metal extrusion, and metal spinning. Study of laminar boundary layer flow caused by a
55 moving rigid surface was initiated by Sakiadis [1] and later the work was extended to the flow
56 due to stretching of a sheet by Crane [2]. The flow of an incompressible fluid past a moving
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4 surface has several engineering applications. The aerodynamic extrusion of plastic sheets, the
5 cooling of a large metallic plate in a cooling bath, the boundary layer along a liquid film in
6 condensation process and a polymer sheet or filament extruded continuously from a die, or a
7 long thread traveling between a feed roll and a wind-up roll are the examples of practical
8 applications of a continuous flat surface. In certain dilute polymer solution (such as 5.4% of
9 polyisobutylene in cetane and 0.83% solution of ammonium alginate in water [3,4]), the
10 viscoelastic fluid flow occurs over a stretching sheet. Any fluid that does not behave in
11 accordance with the Newtonian constitutive relation is called non-Newtonian [5–12]. Non-
12 Newtonian fluids have gained considerable importance because the power required in stretching
13 a sheet in a viscoelastic fluid is less than when it is placed in a Newtonian fluid; and the heat
14 transfer rate for a viscoelastic fluid is found to be less than that of Newtonian fluid.

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16 Recently, Aziz [13] examined a similarity solution for laminar thermal boundary layer over a flat
17 plate with a convective surface boundary condition. Makinde and Olanrewaju [14] and Makinde
18 et al. [15] extended Aziz [13] work by adding the buoyancy effects on thermal boundary layer
19 and internal heat generation term to Aziz [13] work. Other papers that are relevant to the present
20 work are those of Wang [16] and Anderson [17]. The papers by Abel and Mahesha [18, 19] are
21 important contributions because they included effect of variable thermal conductivity, heat
22 source, radiation, buoyancy, magneto-hydrodynamic effects, and viscoelastic behavior of the
23 fluid.

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25 This present work examined the combined effects of internal heat generation and buoyancy
26 effects on thermal boundary layer over a vertical plate with a convective surface boundary
27 condition. Using a similarity approach, the governing equations are transformed into nonlinear
28 ordinary differential equations and solved numerically using a shooting iteration technique
29 together with fourth order Runge-Kutta integration scheme. The pertinent results are displayed
30 graphically and discussed quantitatively

31 32 33 34 35 36 37 38 39 40 41 42 43 44 45 46 47 48 49 50 51 52 **2. Mathematical analysis**

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54 We consider the steady flow of a stream of cold incompressible fluid at temperature T_∞ over the
55 right surface of the vertical plate with a uniform velocity U_∞ while the left surface of the plate is
56 heated by convection from a hot fluid at temperature T_f which provides a heat transfer coefficient
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4 h_f . The cold fluid in contact with the right surface of the plate generates heat internally at the
5 volumetric rate \dot{q} . The density variation due to buoyancy effects is taken into account in the
6 momentum equation using the Boussinesq approximation. The continuity, momentum, and
7 energy equations describing the flow can be written as
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$$10 \quad \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0, \quad (1)$$

$$11 \quad u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = \nu \frac{\partial^2 u}{\partial y^2} + g\beta(T - T_\infty), \quad (2)$$

$$12 \quad \rho c_p (u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y}) = k \frac{\partial^2 T}{\partial y^2} + \dot{q} \quad (3)$$

13 where u and v are the x (along the plate) and the y (normal to the plate) components of the
14 velocities, respectively, T is the local temperature, ν is the kinematics viscosity of the fluid, ρ is
15 the fluid density, c_p is the specific heat at constant pressure and k is the thermal conductivity of
16 the fluid. The velocity boundary conditions can be expressed as
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$$19 \quad u(x, 0) = v(x, 0) = 0, \quad u(x, \infty) = U_\infty. \quad (4)$$

20 The thermal boundary conditions at the plate left surface and far into the cold fluid at the plate
21 right surface may be written as
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$$24 \quad -k \frac{\partial T}{\partial y}(x, 0) = h_f [T_f - T(x, 0)], \quad (5)$$

$$25 \quad T(x, \infty) = T_\infty, \quad (6)$$

26 Introducing a similarity variable η and a dimensionless stream function $F(\eta)$ and temperature
27 $\theta(\eta)$ as
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$$\eta = \frac{y}{x} \sqrt{\text{Re}_x}, u = U_\infty F', v = \frac{\nu}{2x} \sqrt{\text{Re}_x} (\eta F' - F), \theta = \frac{T - T_\infty}{T_f - T_\infty}, \quad (7)$$

$$\lambda_x = \frac{\dot{q} x^2 e^\eta}{k \text{Re}_x (T_f - T_\infty)}, Gr_x = \frac{\nu x g \beta (T_f - T_\infty)}{U_\infty^2},$$

where prime symbol denotes differentiation with respect to η and $\text{Re}_x = U_\infty x / \nu$ is the local Reynolds number. The local internal heat generation parameter λ_x is defined so that the internal heat generation \dot{q} decays exponentially with the similarity variable η as stipulated in [6] and the local thermal Grashof number Gr_x . This type of model can be used in mixtures where a radioactive material is surrounded by inert alloys and in the electromagnetic heating of materials [18]. After substituting Eq.(8) into Eqs. (1) – (6), we obtain the following locally similar equations:

$$F''' + \frac{1}{2} FF' + Gr_x \theta = 0 \quad (8)$$

$$\theta'' + \frac{1}{2} \text{Pr} F \theta' + \lambda_x e^{-\eta} = 0 \quad (9)$$

$$F(0) = F'(0) = 0, \theta'(0) = -Bi_x [1 - \theta(0)] \quad (10)$$

$$F'(\infty) = 1, \theta(\infty) = 0, \quad (11)$$

where

$$Bi_x = \frac{h_f}{k} \sqrt{\frac{\nu x}{U_\infty}}, \text{Pr} = \frac{\rho c_p \nu}{\alpha}. \quad (12)$$

The solution generated whenever Bi_x , λ_x and Gr_x are defined as in Eqs. (8)-(12) are local similarity solutions. In order to have a true similarity solution the parameters Bi_x , λ_x and Gr_x must be constants and not depend on x . This condition can be met if the heat transfer coefficient h_f is proportional to $x^{-1/2}$, the internal heat generation \dot{q} is proportional to x^{-1} and the thermal expansion coefficient β is proportional to x^{-1} . We therefore assume

$$h_f = cx^{-1/2}, \beta = mx^{-1}, \dot{q} = lx^{-1}, \quad (13)$$

where c , m and l are constants but have the appropriate dimensions. Substituting Eq. (13) into Eqs. (7) and (12), we obtain

$$Bi = \frac{c}{k} \sqrt{\frac{\nu}{U_\infty}}, Gr = \frac{\nu m g (T_f - T_\infty)}{U_\infty^2}, \lambda = \frac{l \nu e^\eta}{\alpha U_\infty (T_f - T_\infty)}. \quad (14)$$

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4 The Biot number lumps together the effects of convection resistance of the hot fluid and the
5 conduction resistance of the flat plate. The parameter λ is a measure of the strength of the
6 internal heat generation and the parameter Gr is the thermal Grashof number.
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10 **3. Numerical Solutions**

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12 The coupled nonlinear Eqs (8) and (9) with the boundary conditions in Eqs. (10) and (11) are
13 solved numerically using the fourth-order Runge-Kutta method with a shooting technique and
14 implemented on Maple [20]. The step size is used to obtain the numerical solution with seven-
15 decimal place accuracy as the criterion of convergence.
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22 **4. Results and Discussion**

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24 Figures 2-9 illustrate the influence of the local Biot number Bi_x , local Grashof number Gr_x (the
25 buoyancy effect), local internal heat generation parameter λ_x and the Prandtl number Pr on the
26 velocity $F'(\eta)$, temperature $\theta(\eta)$, the local skin friction coefficient and the local Nusselt number,
27 respectively. Comparison is made with previous results in tables 1 and 2 while table 3 shows the
28 influence of embedded parameters on the overall flow structure. Attention is focused on positive
29 values of the buoyancy parameters i.e. local Grashof number $Gr_x > 0$ (which corresponds to the
30 cooling problem). It is clearly seen in tables 1 and 2 that the special cases of our results are in
31 perfect agreement with those reported in [13-15]. From table 3, we observed that the local skin-
32 friction and the rate of heat transfer at the plate right surface decreases as local Grashof number
33 and local Biot number increases while the local skin-friction and the rate of heat transfer at the
34 plate right surface increases as the internal heating parameter increases. Increase in Prandtl
35 number brings an increase in the local skin-friction and the rate of heat transfer at the plate right
36 surface. Figure 2 depicts the velocity profiles for various values of local Grashof number with
37 other parameters remain constant. It was observed that increase in local Grashof number bring an
38 increase in the velocity which thickens the velocity boundary layer. Moreover, the fluid velocity
39 increases from the plate right surface, attains its peak value within the boundary layer and
40 decreases to its free stream values satisfying the boundary condition. Similar trend is observed in
41 figure 3 with respect to an increase in the internal heat generation parameter. The fluid within the
42 boundary layer becomes lighter and flow faster due to internal heating. Figures 4 and 5 show the
43 effect of increasing local Biot number and Prandtl number. It is interesting to note that the
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velocity boundary layer thickness decreases as local Biot number and Prandtl number increase. Figures 6-9 illustrate the fluid temperature profiles within the boundary layer. The fluid temperature is maximum at the plate right surfaces and decreases exponentially to zero value far away from the plate satisfying the boundary conditions. From these figures, it is noteworthy that the thermal boundary layer thickness increases with a decrease in local Biot number, Prandtl number and local Grashof number while it increases as the local internal heat generation parameter increases.

Table 1: Computations showing comparison with Aziz [13] results for $Gr_x=0$, $\lambda_x=0$ and $Pr=0.72$

Bi_x	Aziz[13] $\theta(0)$	Aziz [13] $-\theta'(0)$	Present $\theta(0)$	Present $-\theta'(0)$
0.05	0.1447	0.0428	0.1447	0.0428
0.10	0.2528	0.0747	0.2528	0.0747
0.20	0.4035	0.1193	0.4035	0.1193
0.40	0.5750	0.1700	0.5750	0.1700
0.60	0.6699	0.1981	0.6699	0.1981
0.80	0.7302	0.2159	0.7302	0.2159
1.00	0.7718	0.2282	0.7718	0.2282
5.00	0.9441	0.2791	0.9441	0.2791
10.00	0.9713	0.2871	0.9713	0.2871
20.00	0.9854	0.2913	0.9854	0.2913

Table 2: Computations showing comparison with Makinde & Olanrewaju [14] results

Bi_x	Pr	λ_x	Makinde& Olanrewaju[14] $\theta'(0)$	Makinde& Olanrewaju[14] $\theta(0)$	Present $\theta'(0)$	Present $\theta(0)$
0.1	0.72	1	0.1154879	2.15487958	0.1154879	2.15487958
1.0	0.72	1	0.3526541	1.35265410	0.3526541	1.35265410
10	0.72	1	0.4437910	1.04437910	0.4437910	1.04437910
0.1	3.0	1	0.0272290	1.27229008	0.0272290	1.27229008
0.1	7.10	1	-0.0101008	0.89899201	-0.0101008	0.89899201
0.1	0.72	5	0.8763365	9.76336572	0.8763365	9.76336572
0.1	0.72	10	1.8273973	19.273973	1.8273973	19.2739733

Table 3: Computation showing $f''(0), \theta'(0)$ and $\theta(0)$ for different parameter values

Bi_x	Pr	λ_x	Gr_x	$f''(0)$	$\theta'(0)$	$\theta(0)$
0.1	0.72	1	0.1	0.639970	0.089403	1.894034
1.0	0.72	1	0.1	0.570644	0.303644	1.303644
10	0.72	1	0.1	0.538641	0.396797	1.039679
0.1	3.0	1	0.1	0.479768	0.016571	1.165719
0.1	7.10	1	0.1	0.421585	-0.01547	0.845266
0.1	0.72	5	0.1	1.319375	0.611718	7.1171820
0.1	0.72	10	0.1	1.916164	1.157887	12.578876
0.1	0.72	1	1.0	2.023402	0.040122	1.4012239
0.1	0.72	1	10	7.698598	-0.01036	0.8963780
0.1	0.72	1	20	11.564248	-0.023012	0.769870

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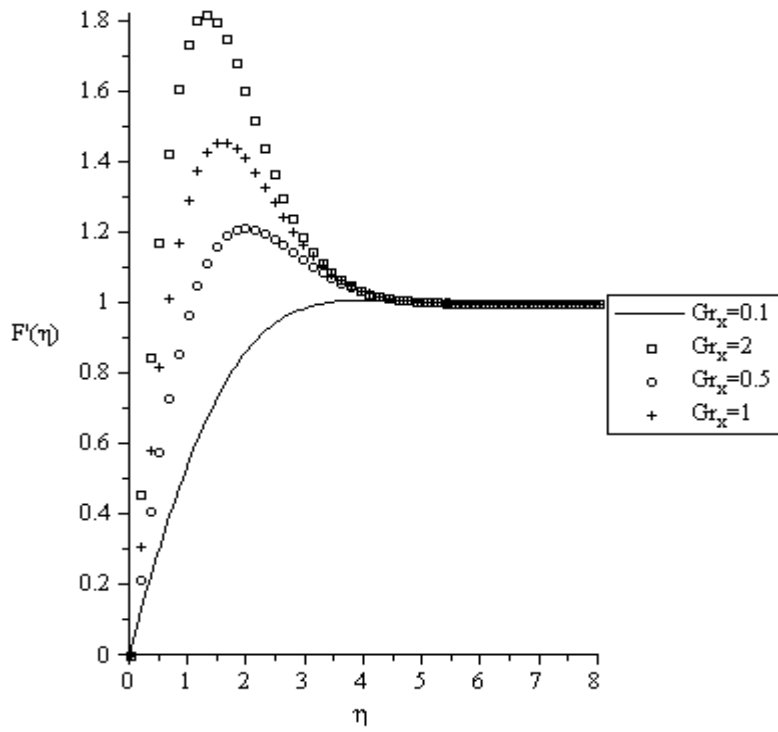


Figure 2: Velocity profiles for $Bi=0.1, \lambda_x=1, Pr=0.72$

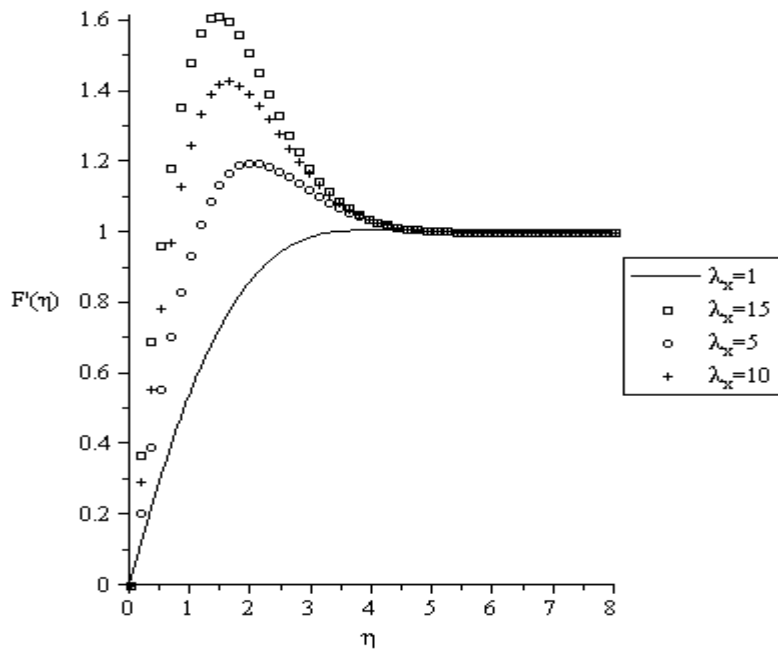


Figure 3: Velocity profiles for $Gr_x=0.1, Bi=0.1, Pr=0.72$

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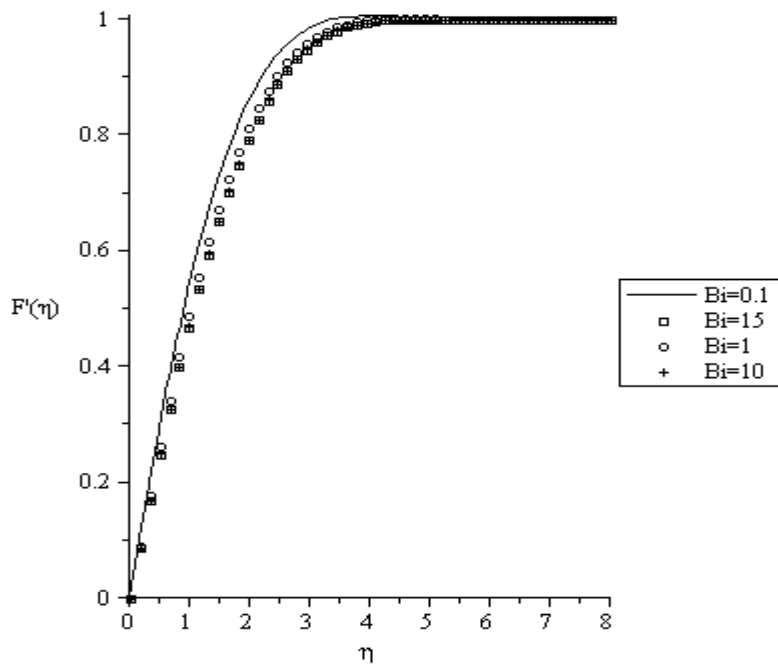


Figure 4: Velocity profiles for $Gr_x=0.1$, $\lambda_x=1$, $Pr=0.72$

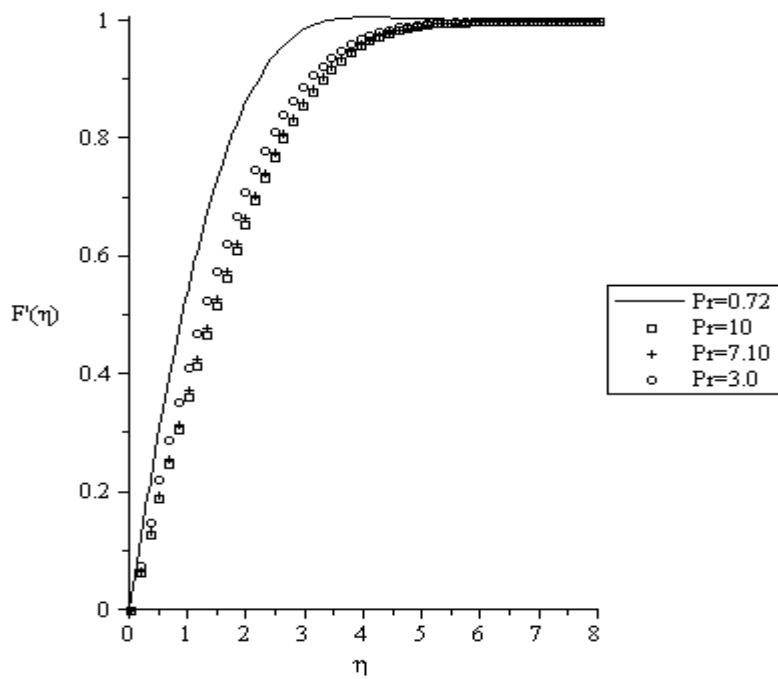


Figure 5: Velocity profiles for $Gr_x=0.1$, $Bi=0.1$, $\lambda_x=1$

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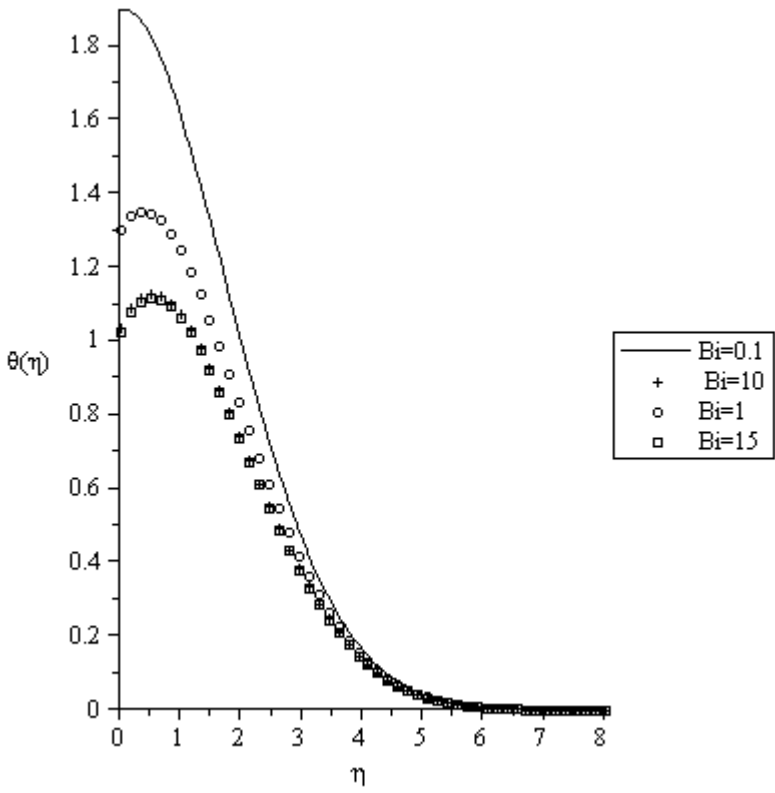


Figure 6: Temperature profiles for $Gr_x=0.1$, $\lambda_x=1$, $Pr=0.72$

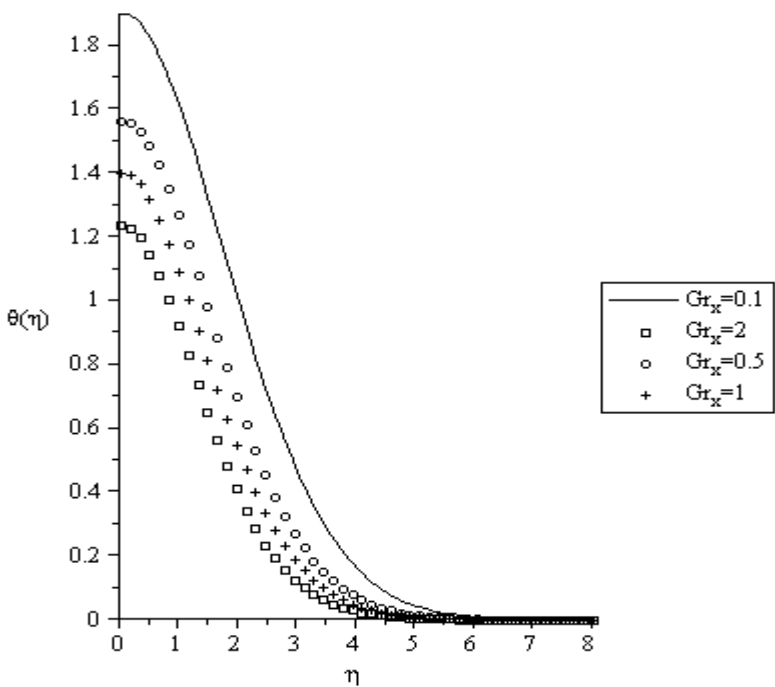


Figure 7: Temperature profiles for $Bi=0.1$, $\lambda_x=1$, $Pr=0.72$

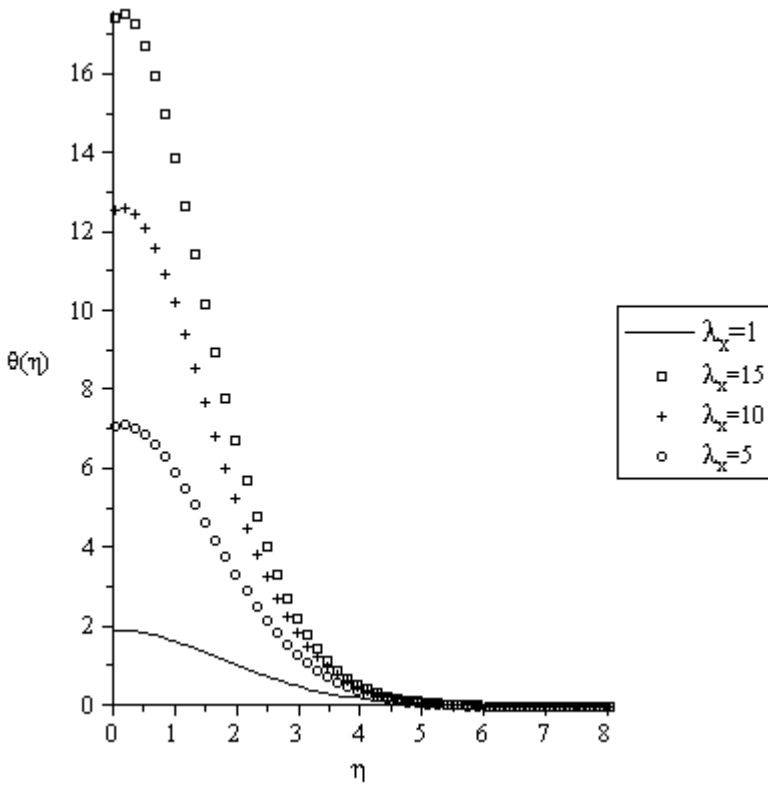


Figure 8: Temperature profiles for $Gr_x=0.1$, $Bi=0.1$, $Pr=0.72$

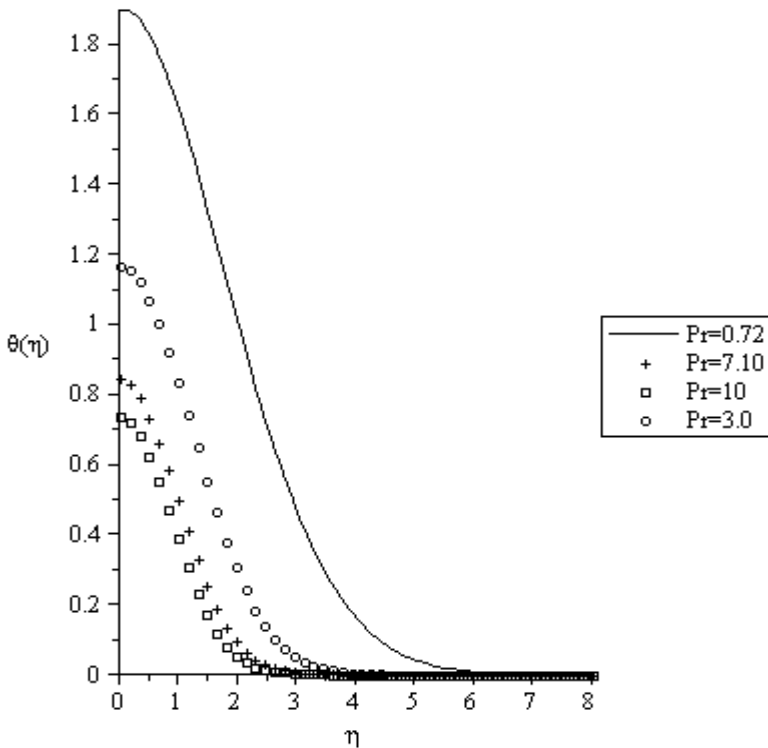


Figure 9: Temperature profiles for $Gr_x=0.1$, $Bi=0.1$, $\lambda_x=1$

5. Conclusions

The problem on boundary layer flow past a vertical plate due to gravity and fluid density variation due to temperature with internal heat generation and buoyancy effects has been considered. Using similarity variable and the fourth-order Runge-Kutta method coupled with shooting technique, the governing equations are tackled numerically and the influence of various embedded parameters have been discussed quantitatively. Our results reveal among others that;

- the velocity boundary layer thickness increases with an increase in local Grashof number due to buoyancy effects and local internal heat generation.
- the thermal boundary layer thickness increases with a decrease in local Biot number, Prandtl number and local Grashof number while it increases as the local internal heat generation parameter increases.
- the local skin-friction and the rate of heat transfer at the plate right surface decreases as local Grashof number and local Biot number increases while the local skin-friction and the rate of heat transfer at the plate right surface increases as the internal heating parameter increases.

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